

Equations of motion. For example:

$$(\partial_t + \mathbf{u} \cdot \nabla) \mathbf{u} + 2\boldsymbol{\Omega} \times \mathbf{u} + \frac{\nabla p}{\rho} - \mathbf{g} = \nu \nabla^2 \mathbf{u}$$

$$(\partial_t + \mathbf{u} \cdot \nabla) \theta = \kappa \nabla^2 \theta$$

$$(\partial_t + \mathbf{u} \cdot \nabla) S = \mu \nabla^2 S$$

and bndry / initial conditions. This is *really old stuff*.

Our job is numeric approx solutions (“modeling”)

Equations of motion. For example:

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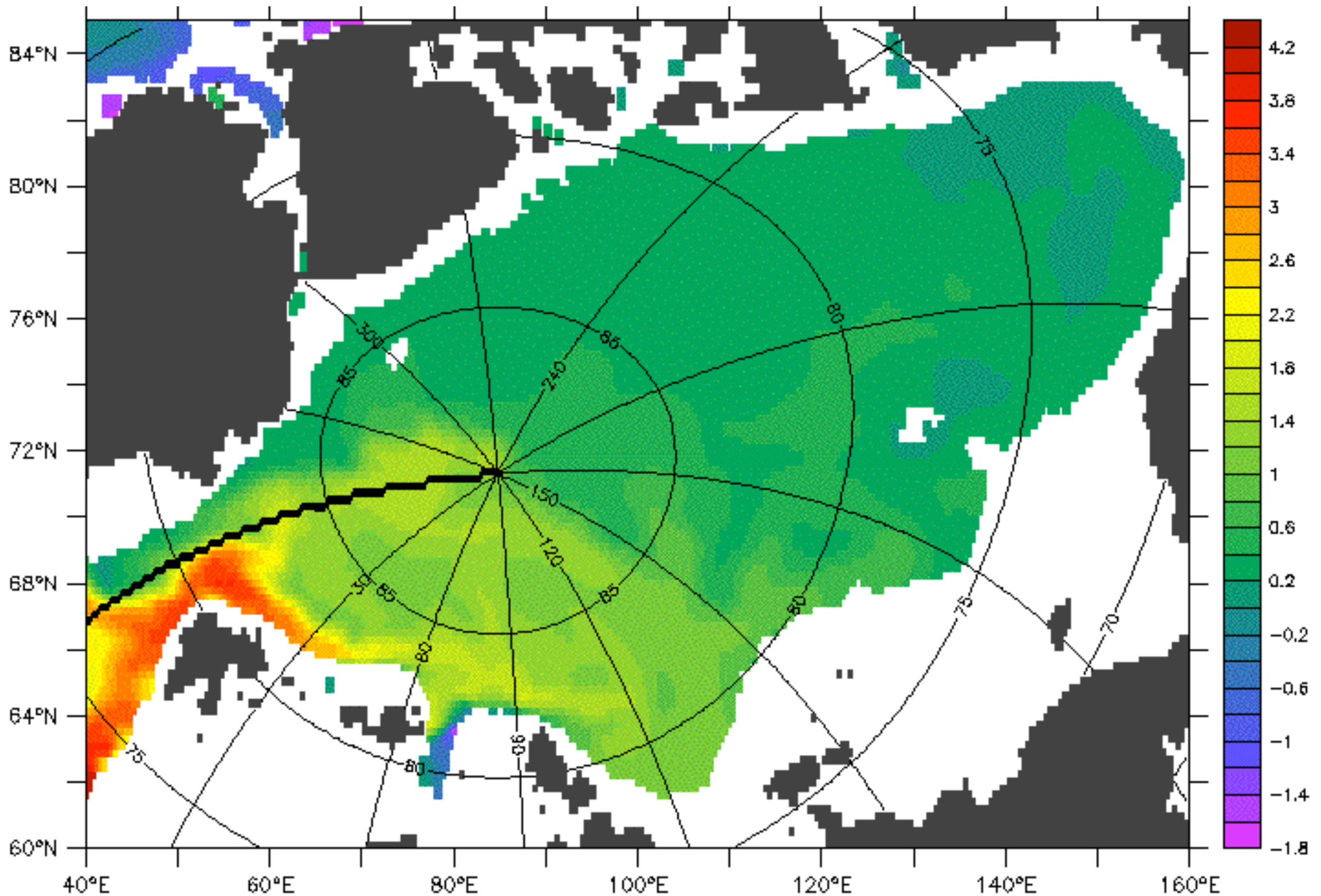
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Our job is numeric approx solutions (“modeling”)
-- maybe fudging some coefs (“parameterization”).

This talk: Rethink equations of motion. Why?
Consider dependent variables (stuff that “moves”)?

Hint: it matters.

What is this stuff? θ at some x,y,z,t ? How about $\langle \theta \rangle$?



Rewrite symbolically $\frac{\partial \mathbf{Y}}{\partial t} = \mathbf{F}(\mathbf{Y}) + \mathbf{G}$

where \mathbf{Y} collects everything that interests you.

This equation, solved honestly, requires mm grid.
Ignore small scale details \Rightarrow large scale expectations.

Probability! $dP = P(\mathbf{Y})d\mathbf{Y}$ $\langle \mathbf{Y} \rangle \equiv \int \mathbf{Y} dP$

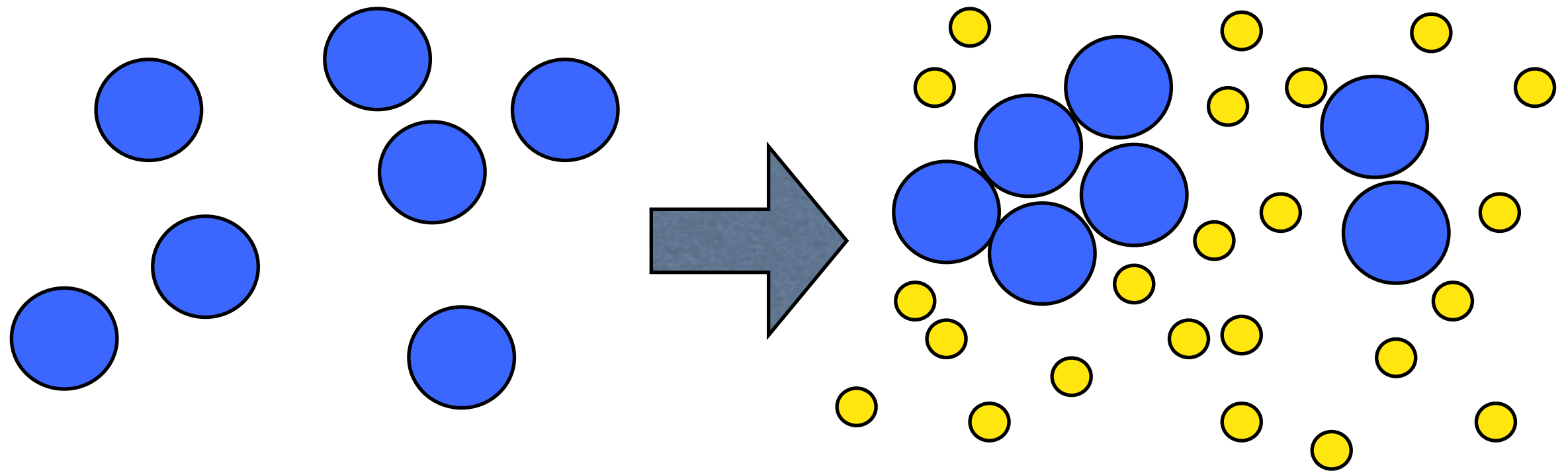
$\frac{\partial \langle \mathbf{Y} \rangle}{\partial t} = \langle \mathbf{F}(\langle \mathbf{Y} \rangle) \rangle + \langle \mathbf{G} \rangle + \text{"oops"}$ *What to do about "oops"?*

Entropy:

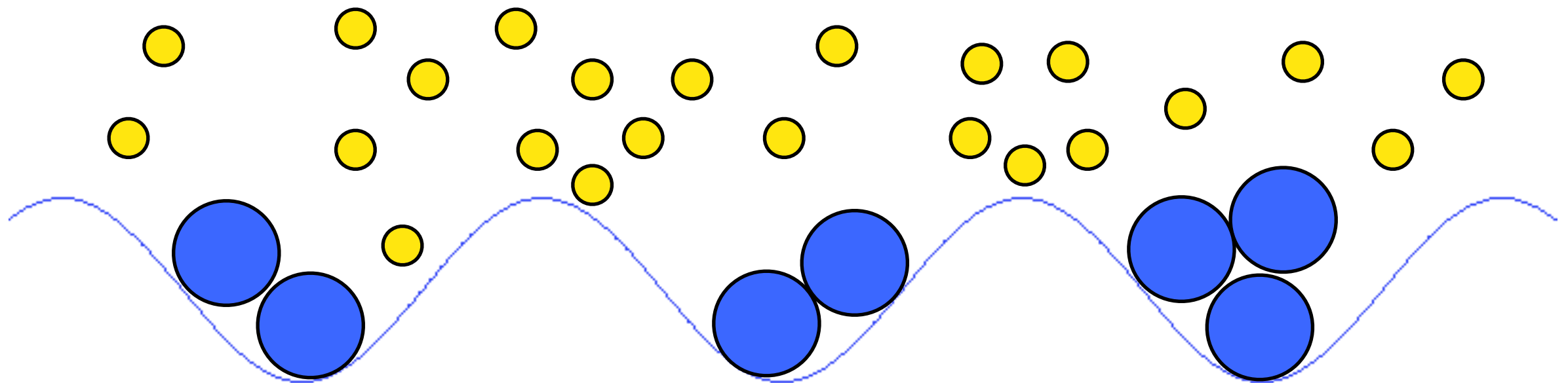
$$H \equiv - \int dP \log P$$

then $\frac{\partial \langle \mathbf{Y} \rangle}{\partial t} \stackrel{?}{\approx} \langle \mathbf{F}(\langle \mathbf{Y} \rangle) \rangle + \langle \mathbf{G} \rangle + \mathbf{C} \cdot \partial_{\langle \mathbf{Y} \rangle} H$

entropic force



Examples from nanoworld (colloids, ‘machines’, microbot): The only explicit physics is repulsion among balls, and from walls. “See” attraction. “Entropic forcing” in the lab!



Three illustrations: neptune
 vertical viscosity
 sea ice

neptune (level 1): $\frac{\partial \mathbf{Y}}{\partial t} = \mathbf{F}(\mathbf{Y}) + \mathbf{G} + \mathbf{C} \cdot \partial_{\mathbf{Y}} H$

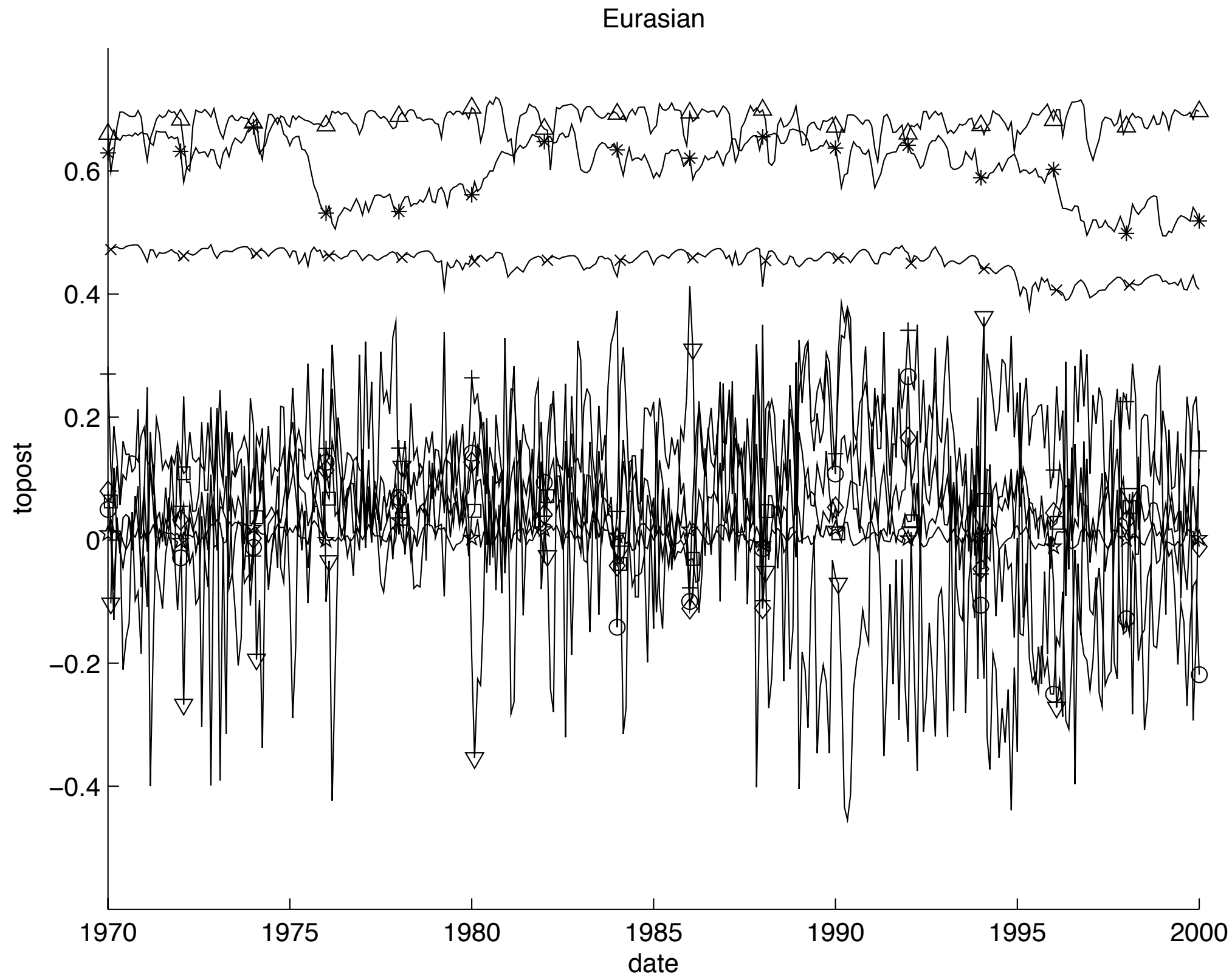
$$\mathbf{C} \cdot \partial_{\mathbf{Y}} H \approx \mathbf{C} \cdot \partial_{\mathbf{Y}} \partial_{\mathbf{Y}} H \cdot (\mathbf{Y}^* - \mathbf{Y}) \equiv \mathbf{K} \cdot (\mathbf{Y}^* - \mathbf{Y})$$

$$\mathbf{u}^* = -L^2 \mathbf{f} \times \nabla \log D \quad (\text{QG energy \& enstrophy})$$

Replace $\partial_t \mathbf{u} + \dots = \partial_z A_z \partial_z \mathbf{u} + \nabla_h A_h \nabla_h \mathbf{u} + \dots$

with $\partial_t \mathbf{u} + \dots = \partial_z A_z \partial_z \mathbf{u} + \underline{\nabla_h A_h \nabla_h (\mathbf{u} - \mathbf{u}^*)} + \dots$

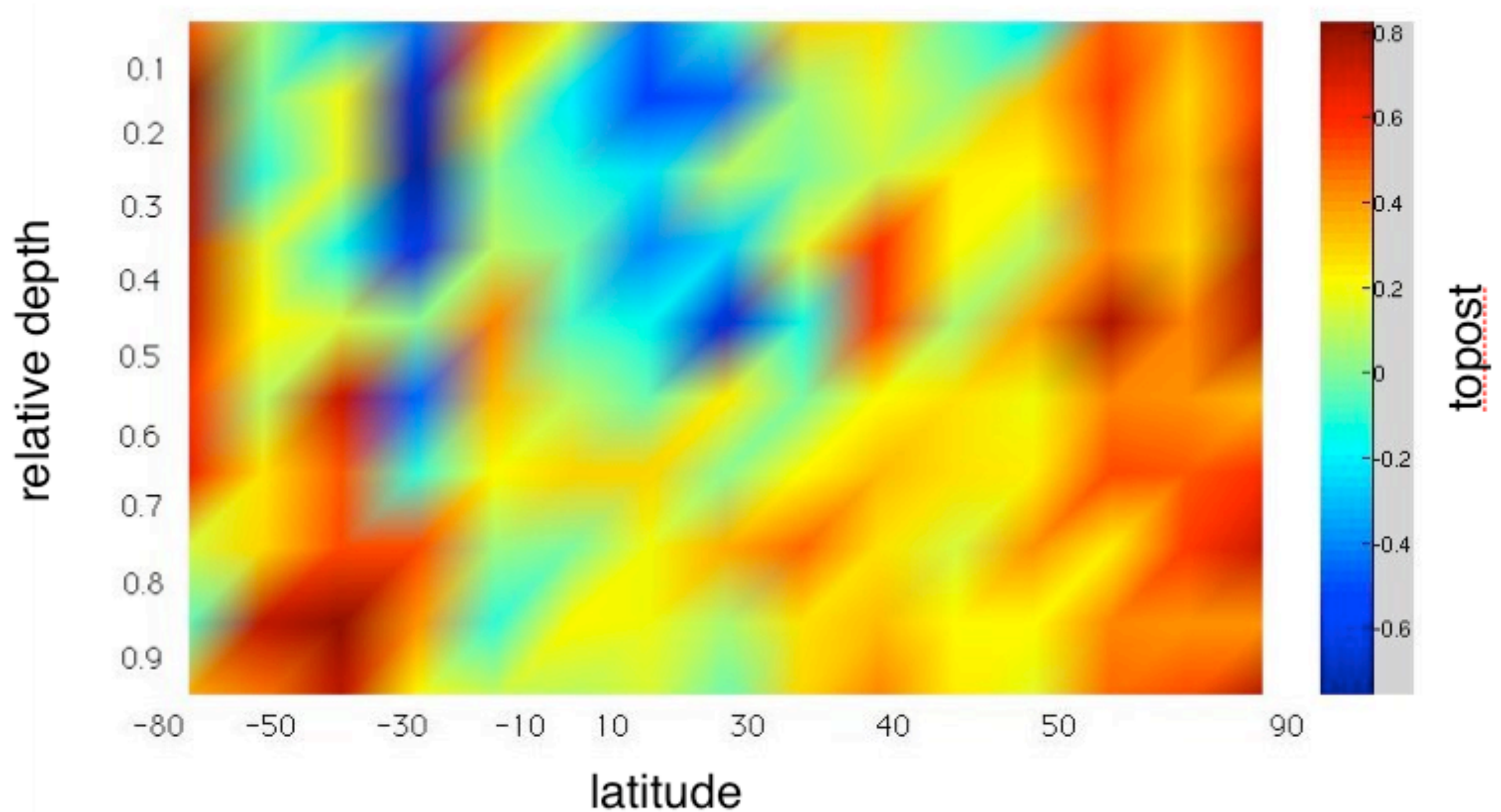
Topostrophy: $\mathbf{f} \times \mathbf{u} \cdot \nabla D / \sqrt{|\mathbf{f} \times \mathbf{u}|^2 |\nabla D|^2}$, basin - averaged from 10 models (JGR, 2007)



Neptune makes big difference but ... **better or worse?**

Observed from 17120 current meters over 83087 months (JGR, 2008):

Topostrophy vs. latitude and relative depth



Observations are completely outside the range of traditional modeling????

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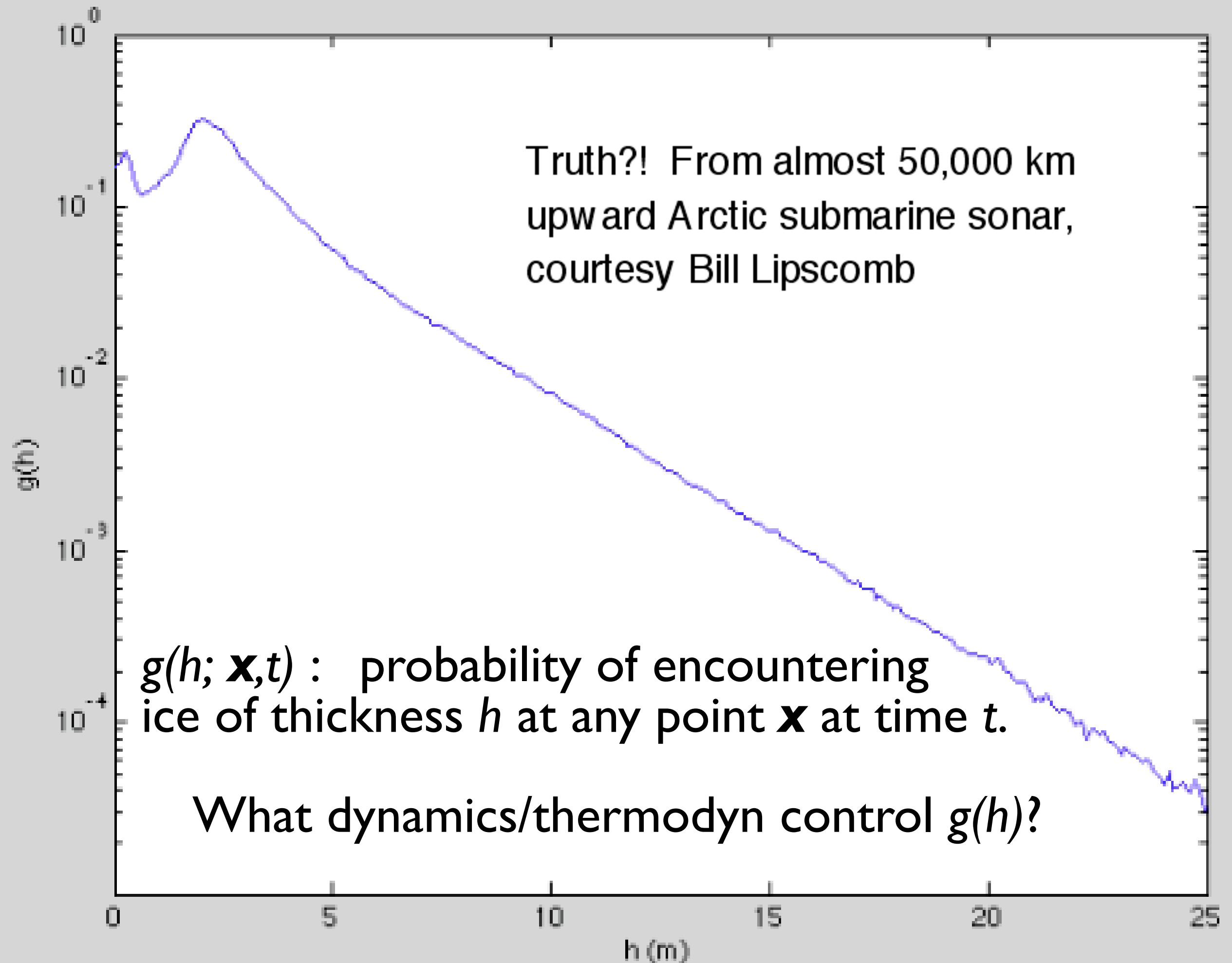
with $\partial_t \mathbf{u} + \dots = \partial_z A_z \partial_z \mathbf{u} + \nabla_h A_h \nabla_h (\mathbf{u} - \mathbf{u}^*) + \dots$

vertical viscosity: $A_z = (f^2 / N^2) A_\eta$

$$f^2 / N^2 \approx 10^{-3}, \quad A_\eta \approx 10^2, \quad \text{then } \underline{A_z \approx 10^{-1} \text{ m}^2 / \text{s}} \quad (\underline{\text{not } 10^{-4} \text{ to } 10^{-3}})$$

Three illustrations: neptune
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A (very!) little math:

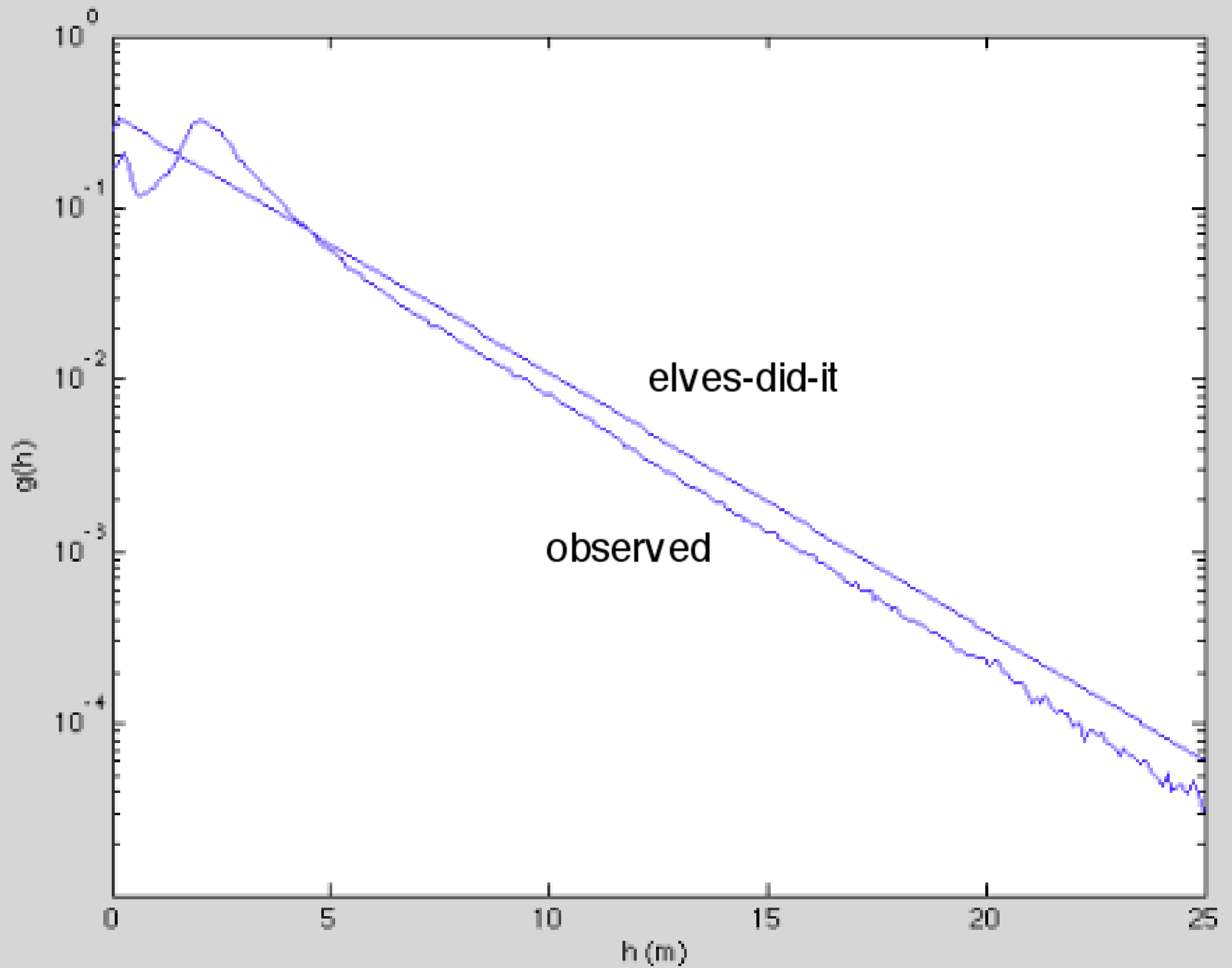
Maximise $S = -\int_0^{\infty} \log(g) g dh$

subject to $\int_0^{\infty} g dh = 1$ and $\int_0^{\infty} g h dh = H$

and *presto!* $g(h) = \exp(-h/H) / H$

Thus, Santa's elves, underoccupied during off-season, randomly toss ice into the Arctic. Is that easy, or what?

Let's see!



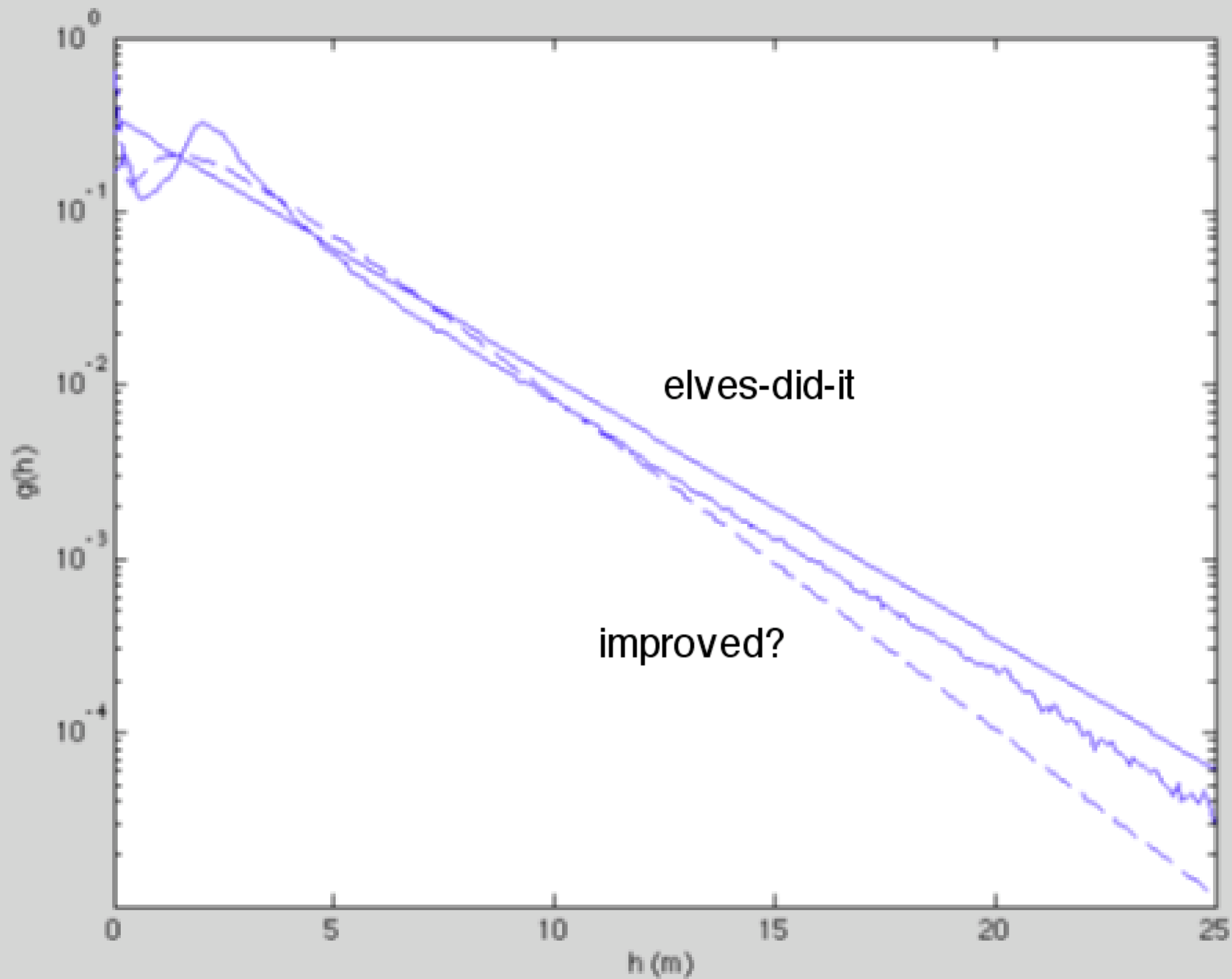
OK, not the greatest ever. But fiddle just a little.

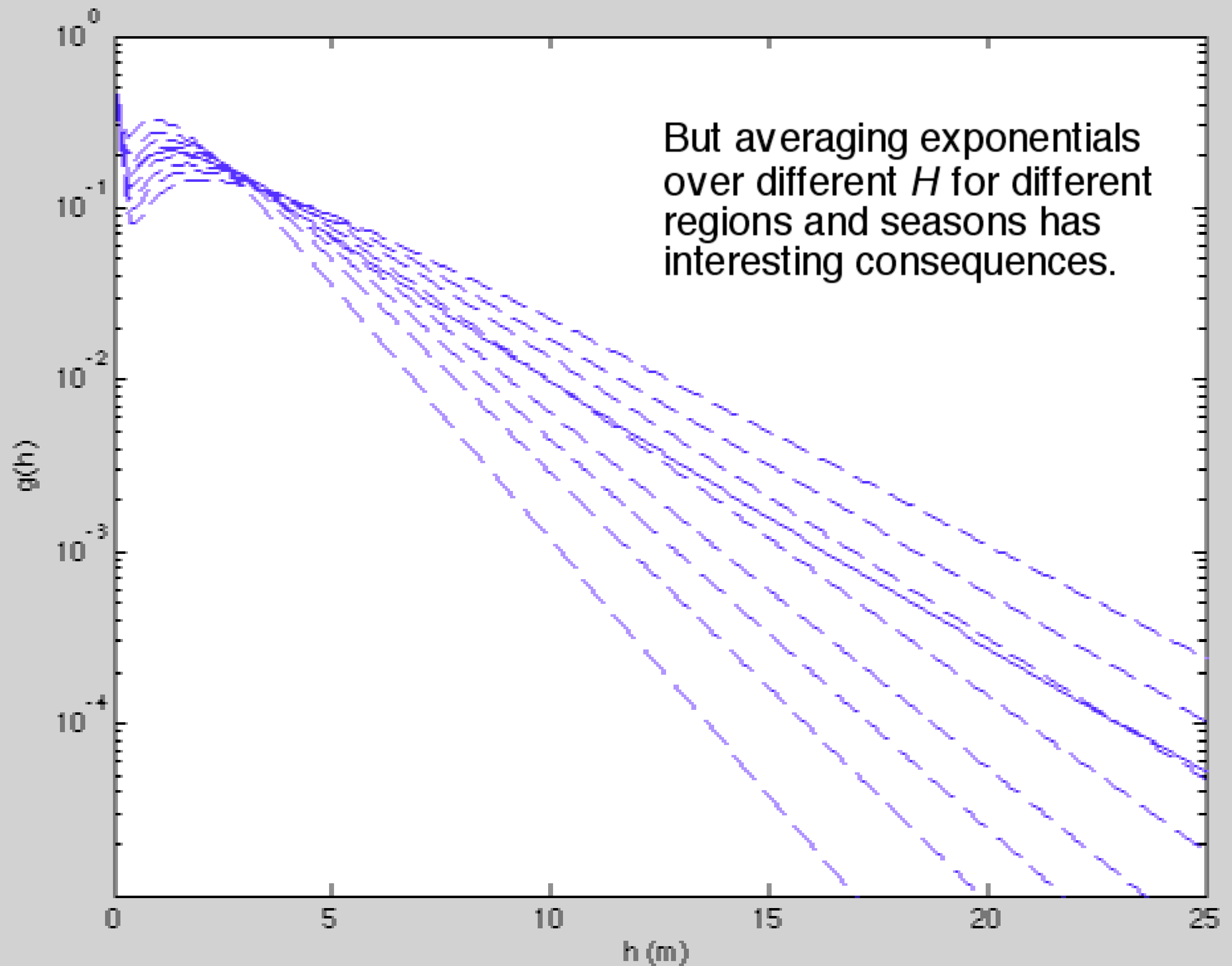
Wind shifts, tides, inertial oscillations, ... open ice.

Recognize an “open water” (thin ice?) fraction $1-A$.
In leads during freezing, thin new ice rapidly forms
but is easily smunched into thicker ice. Fix it?

$$D = \frac{A}{(2a-1)H} \left(e^{-\frac{h}{aH}} - e^{-\frac{h}{H-aH}} \right) + \frac{1-A}{bH} e^{-\frac{h}{bH}}$$

Improved? Let's see.





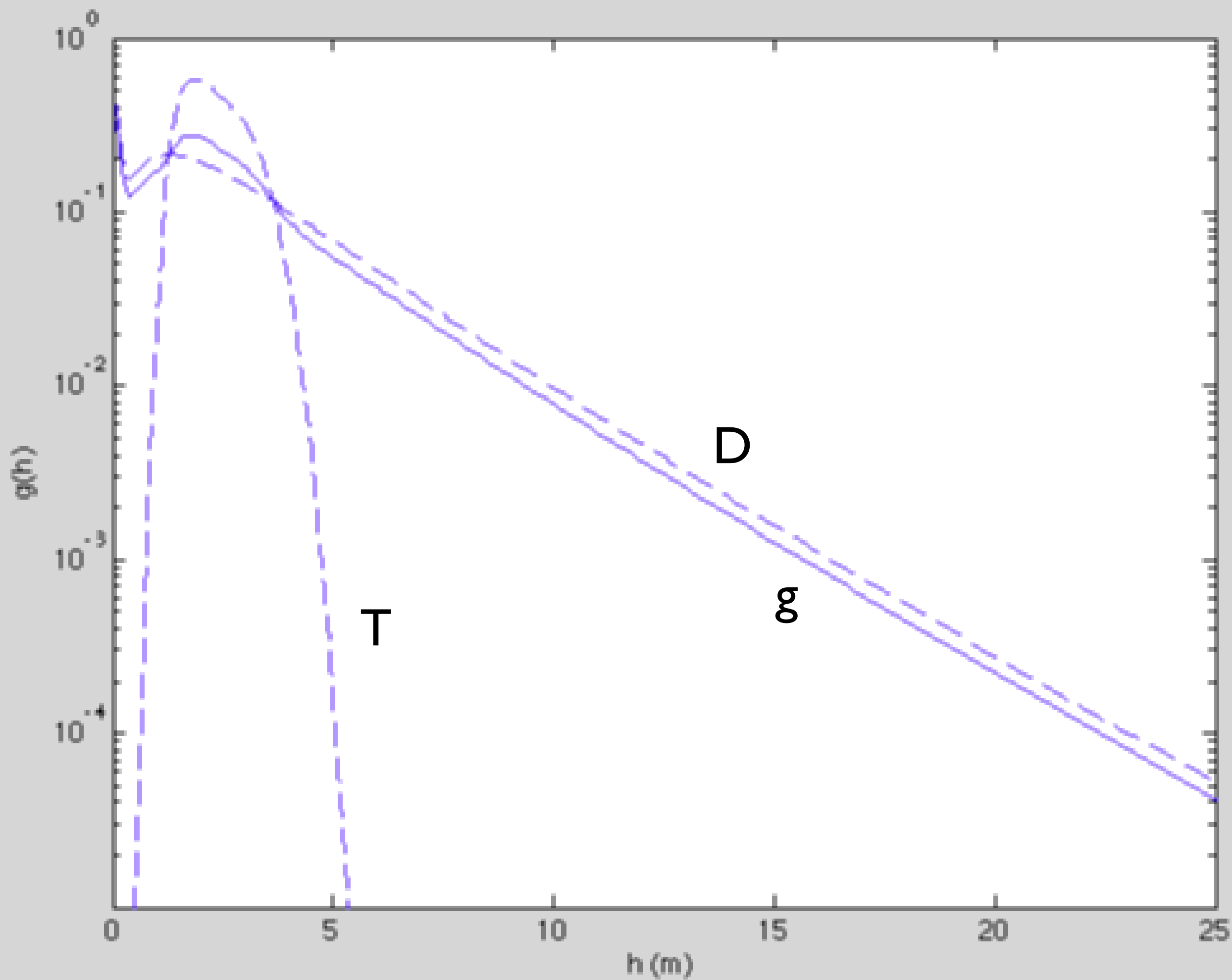
Dynamics -- alone -- spreads out $g(h)$, approaching D .

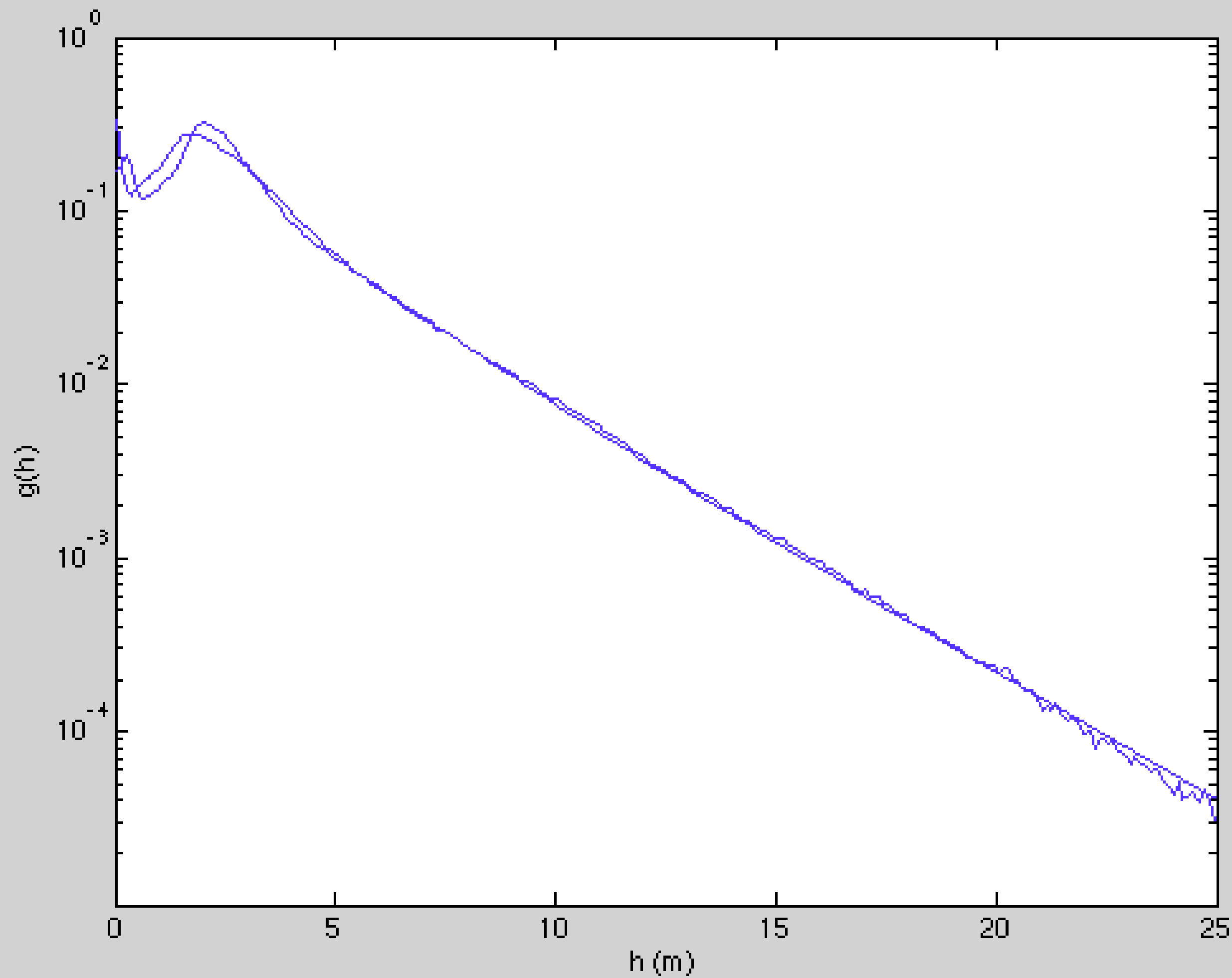
Thermodynamics -- alone -- would focus $g(h)$, e.g. as

$$T = e^{-\frac{(h-aH)^2}{2b^2}} / \sqrt{2\pi b} \quad \text{which also needs be averaged over}$$

various H . Combine to $g(h) = cD + (1-c)T$ for some c .

See a case $c=0.8$:





This is just fudge factors!

Yes. But not so many. b only need be small, like $b \sim 0.1$. a is constrained, $0.5 < a < 1$.
I chose $a = 0.7$. That leaves c . I chose $c = 0.8$
(more dynamic, not so thermodyn.)

Models solve dA/dt and dH/dt . Fudge something for dc/dt , like integrated deformation? Keep first- and multi-year fractions?

Summary:

Stat mech completes the equations of motion. Includes entropic forcing.

E.g., neptune

E.g., vertical viscosity

Entropy also is about economy (intellectual & computational). When outcomes depend upon many detailed interactions, problems get easier. Entropy calculus (a crutch) *avoids* difficult stuff.

E.g., sea ice thickness